

## ARTICLE

DOI: 10.1038/s41467-018-03752-5 **OPEN**

# Rolling up transition metal dichalcogenide nanoscrolls via one drop of ethanol

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Two-dimensional transition metal dichalcogenides (TMDs) have attracted lots of interest because of their potential for electronic and optoelectronic applications. Atomically thin TMD flakes were believed capable to scroll into nanoscrolls (NSs) with distinct properties. However, limited by mechanical strength and chemical stability, production of high-quality TMD NSs remains challenging. Here, we scroll chemical vapor deposition-grown monolayer TMD flakes into high-quality NSs in situ in 5 s with a nearly 100% yield by only one droplet of ethanol solution. An obvious photoluminescence is demonstrated in NSs and the self-encapsulated structure makes NSs more insensitive to external factors in optical and electrical properties. Furthermore, based on the internal open topology, NSs hybridized with a variety of functional materials have been fabricated, which is expected to confer TMD NSs with additional properties and functions attractive for potential application.

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tomically two-dimensional (2D) transition metal dichal-<br>cogenides (TMDs), such as  $MoS<sub>2</sub>$ ,  $Mose<sub>2</sub>$ ,  $WS<sub>2</sub>$ , and  $WSe<sub>2</sub>$ <br>display numerous exceptional electronic and optical<br>properties arising from quantum c cogenides (TMDs), such as  $MoS<sub>2</sub>$ , MoSe<sub>2</sub>, WS<sub>2</sub>, and WSe<sub>2</sub>, properties arising from quantum confinement $1,2$ . Complementary to gapless graphene, 2D TMDs with intrinsic band gaps are promising interesting field-effect transistor (FET) and optoelectronic devices. Recently, the main focus of attention has been for the production of their intrinsic or heterojunction structures, their properties, and their applications on  $2D$  scale<sup>3–8</sup>. In addition to changes in 2D size and form, self-assembly of the atomically thin TMD flakes, as an emerging area, is largely unexplored. Assembly processes by folding and scrolling, i.e., rolling up, can transform relatively simple structures into complex topologies, such as nanoscrolls (NSs), with distinct properties as well as the original excellent characteristics. Indeed, theoretical calculations have predicted the unique topology of 2D material-based NSs to yield unusual electronic and optical properties $9-12$  $9-12$  $9-12$ ; thus these NSs have promise as building blocks in flexible electronics, microfluidics, energy storage, self-propelled micromachines, and optical resonators<sup>13–18</sup>. However, experimental realization of such high-quality NSs has only been achieved for graphene<sup>[13](#page-6-0)–[16](#page-6-0)</sup> and boron nitride<sup>19</sup>, which exhibited high strength and chemical inertness. Limitations in mechanical strength and chemical stability present difficulties in producing high-quality  $TMD-NSs^{20}$  $TMD-NSs^{20}$  $TMD-NSs^{20}$ . Recently, an argon plasma-assisted method has been demonstrated for the fabrication of amorphous  $MoS<sub>2</sub>-NSs$ , but in this method  $MoS<sub>2</sub>$  sheets failed to scroll effectively, only edge region of the sheet curved while the large central area remains plane. Moreover, upon argon plasma etching, a serious damage has been brought to the original sheets, where almost half of sulfur atoms were removed. The obtained NSs exhibited low crystallinity even to be amorphous<sup>21</sup>. A highquality TMD-NS is highly desirable for both fundamental studies and potential applications; however, its reliable experimental fabrication still remains challenging.

Here, we show a very simple method for fabricating highquality TMD-NSs, which only requires one droplet of ethanol solution to scroll chemical vapor deposition (CVD)-grown monolayer TMD flakes in situ in 5 s with a nearly 100% yield. TMD-NSs are promising optoelectronic materials with potential applications in optoelectronic devices because of the high FET mobility. Owing to self-encapsulated structure, the optical and electrical properties of NSs are more insensitive to external factors, while the electrical performance of 2D materials flakes varies with both the underlying substrate and the environment. In addition, because of the internal open topology, the interlayer spacing of TMD-NSs can be easily expanded to accommodate a variety of functional materials, including organic small molecules, polymers, nanoparticles, and 2D materials, as well as biological substances. These features are very attractive for applications in solar cells, photodetectors, flexible logic circuits, energy storage, and sensors.

### Results

Rolling up CVD-grown TMD monolayer flakes. In a typical experiment, large-area, monolayer TMD flakes were synthesized on a  $SiO<sub>2</sub>/Si$  substrate by CVD. Then, the TMD flakes were scrolled into NSs in just  $\frac{1}{5}$  s with a nearly 100% yield by placing one drop of ethanol solution (volume ratio of ethanol:water = 2:1) on their surface. Other aqueous solutions, such as methanol, tetrahydrofuran, dimethylformamide, and N-methyl-2-pyrrolidone, can also be used for TMD flakes scrolling (Supplementary Fig. 1, 2 and Supplementary Note 1). In addition, the scrolling proceeded well for TMD flakes grown on other substrates, such as  $Si<sub>3</sub>N<sub>4</sub>$  and sapphire (Supplementary Fig. 3). To clarify the TMD-NSs formation process, the entire procedure was recorded using a CCD camera equipped on an optical-microscope (Supplementary Movie 1). A series of steps that might have led to the formation of the TMD-NSs are proposed in Fig. 1. Monolayer TMD flakes were grown on a substrate at a high temperature ( $\geq 720$  °C). During the process of cooling down to room temperature, a strain was introduced to the TMD flakes because of the mismatch in thermal expansion coefficients between the TMD flake and the substrate. Stable TMD flakes were achieved when the strain and substrate adhesion were balanced. When an ethanol solution spread onto the surface of TMD flakes, the liquid film intercalated into the TMD flakes and the substrate, as demonstrated by the contrast color change in the flakes (Supplementary Fig. 4). With the liquid intercalation, part of the TMD flake was first released from the substrate to become freestanding as shown in Fig. 1. Under the built-in strain, the released part of the flake curved out of the plane and continued to roll up to a complete NS in the solution. An energy analysis of the  $MoS<sub>2</sub>$  flakes before and after scrolling was presented in the Supplementary Note 2. The energy barrier required to bend the freestanding flake out of plane was proved to be relatively low, indicating the easy nucleation of a scroll. In addition, the strain energy induced from the thermal mismatch of the  $MoS<sub>2</sub>$  flake was found high enough to activate the formation of scrolls after the liquid intercalation and the adhesion energy of the  $MoS<sub>2</sub>$  flake with substrate decreasing (Supplementary Fig. 5). Strain as the stimulus to the NSs



Fig. 1 Schematic of self-scrolling of chemical vapor deposition (CVD)-based transition metal dichalcogenide (TMD) monolayer flakes. A proposed series of steps leading to the formation of TMD-NSs is shown: (I) initial TMD monolayer growth on a substrate at a high temperature (HT) (e.g., 720, 825 °C); (II) the generation of a strain, balanced with the adhesion from substrate, in the TMD monolayer while cooling down to room temperature (RT); (III) the insertion of liquid between the TMD monolayer and substrate and the disappearance of the adhesion; (IV) the TMD flakes curve out of the plane driven by the strain; (V) the final formation of TMD-NSs

formation was also suggested by the following fact. When treated with ethanol solution, freshly produced TMD flakes on substrates were more prone to scroll in the shortest time with the highest yield, while the aged samples exhibited slightly decreased scrolling speed and yield. A few cracks were observed in aged samples by scanning electron microscopy (SEM, Supplementary Fig. 6), and these cracks could have partially relieved the strain.

TMD flakes were observed to prefer to scroll along one edge where the adhesion with the substrate was first relaxed upon the intercalation of the liquid. As the edges of CVD-grown TMD single crystals have a zigzag orientation<sup>4,22,23</sup>, the chirality of their NSs are defined (Supplementary Fig. 7). For large-area polycrystalline films without regular edges, the orientation and chirality of their NSs are generally random (Fig. 2a, b). Thus, the controllable scrolling of TMD films was attempted (Fig. 2c–e and Supplementary Fig. 8). Using a focused ion beam (FIB), a  $MoS<sub>2</sub>$  polycrystalline film was patterned into parallel ribbons with defined width and direction (Fig. 2c) to produce long NS arrays (Fig. 2d). Using FIB etching a second time, long NSs were further arranged into periodic arrays with a controlled length (Fig. 2e). These NS arrays have promise for wide application, such as integrated circuits and matrix displays. We assume that horizontal TMD-NS arrays with controllable chirality could potentially be achieved by controlling the crystal orientation of the ribbon edges via etching from one large-area TMD single crystal, which is prevented by the technique to synthesize largearea TMD single crystals.

Morphology characterization. CVD-grown TMD flakes are high quality even comparable to intrinsic flakes $23,24$ . The ethanol solution was clean and harmless; therefore, no defects or impurities were introduced into the high-quality TMD-NSs (Supplementary Fig. 9). Based on this method, several typical TMD-NSs, including those of  $MoS<sub>2</sub>$  (Fig. 2f and Supplementary Fig. 10),  $WS<sub>2</sub>$ (Fig. 2g),  $Mose_2$  (Fig. 2h), and  $WSe_2$  (Fig. 2i), were successfully fabricated. Most of the NSs displayed a straight and compact appearance. High-resolution transmission electron microscopy (HR-TEM) was used to further examine the TMD-NSs microstructure, which was found to be multiwalled and tubular (Fig. 2j–m). The diameter and sidewall thickness of the TMD-NS varied with the size of the original flake. The hollow diameters were 10–40 nm. In addition, the high-magnification images of the sidewalls and the selected area electron diffraction (SAED) both showed that the TMD layers in the NSs were stacked uniformly and compactly (Fig. 2j–m inset and Supplementary Fig. 11). The distance between adjacent walls was approximately 6.21 Å, 6.41 Å, 6.35 Å, and 6.45 Å for  $MoS_2$ -,  $WS_2$ -,  $MoSe_2$ -, and  $WSe_2$ -NSs, respectively, which matched the interlayer spacing in bulk TMDs (6.15 Å, 6.36 Å, 6.45 Å, and 6.49 Å for  $MoS_2$ ,  $WS_2$ ,  $MoSe_2$ , and



Fig. 2 Transition metal dichalcogenide (TMD)-nanoscrolls (NSs) from self-scrolling chemical vapor deposition (CVD)-based TMD monolayer flakes. a Optical image of CVD-grown MoS<sub>2</sub> monolayer flakes on a SiO<sub>2</sub>/Si substrate (the purple area represents the substrate, the green area represents the MoS<sub>2</sub> monolayers; scale bar, 500 μm). b Optical image of MoS<sub>2</sub>-NSs on a SiO<sub>2</sub>/Si substrate (scale bar, 100 μm). c-e The fabrication process of a MoS<sub>2</sub>-NS array (scanning electron microscopy—SEM image shown, scale bars, 50  $\mu$ m). c Large-area MoS<sub>2</sub> monolayer film patterned into ribbons by focused ion beam (FIB) etching (the dark region represents MoS<sub>2</sub> film, the bright area is bare substrate).  $d$  Long MoS<sub>2</sub>-NSs made from the controllable scrolling of the patterned MoS<sub>2</sub> film in c. (The white parallel lines are MoS<sub>2</sub>-NSs, and the dark area represents the substrate). e A 12 × 6 array of MoS<sub>2</sub>-NSs fabricated via a second FIB etching of the long MoS<sub>2</sub>-NSs in d. f-i SEM images of typical TMD-NSs on SiO<sub>2</sub>/Si substrates: MoS<sub>2</sub>-NSs (f), WS<sub>2</sub>-NSs (g), MoSe<sub>2</sub>-NSs (h), and WSe<sub>2</sub>-NSs (i). (Scale bars, 5 μm in f, i and 10 μm in g, h). j–m TEM images of typical TMD-NSs: MoS<sub>2</sub>-NSs (j),WS<sub>2</sub>-NSs (k), MoSe<sub>2</sub>-NSs (l), and WSe<sub>2</sub>-NSs (m). (Scale bars, 20 nm). Inset: High-magnification images of sidewalls of TMD-NSs (scale bars, 2 nm)

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Fig. 3 Optical characterization of MoS<sub>2</sub>-NSs (nanoscrolls). a Raman spectra of chemical vapor deposition (CVD)-grown MoS<sub>2</sub> monolayers (MoS<sub>2</sub>-ML, black line) and MoS<sub>2</sub>-NSs (red line) with 532 nm excitation. **b** Optical image of a typical MoS<sub>2</sub>-NS. c Corresponding fluorescence image of the MoS<sub>2</sub>-NS in b. The fluorescence image was obtained with an excitation wavelength of 510-560 nm. d PL spectra from MoS<sub>2</sub> monolayer (black line) and MoS<sub>2</sub>–NSs (red line) (532 nm excitation). e, f PL spectra response of MoS<sub>2</sub> monolayer (e) and MoS<sub>2</sub>-NSs (f) with NH<sub>3</sub> doping for 0 s (black line), 10 s (red line), and 10 min (blue line) (scale bars,  $5 \mu m$  in **b** and **c**)

WSe<sub>2</sub>, respectively), indicating both the tight scrolling of the TMD flakes and the lack of interlayer contamination.

Optical characterization. Raman spectra of the TMD-NSs were recorded using 532 nm excitation (Fig. 3a and Supplementary Fig. 12a–c). The Raman signature of monolayer  $MoS<sub>2</sub>$  flakes is consistent with that of mechanically exfoliated monolayer  $\text{MoS}_2^{25}$ . A similar Raman signature was also found for  $MoS<sub>2</sub>-NSs$ , demonstrating the good crystallinity of the MoS<sub>2</sub>-NSs. A blue shift was observed in the  $A_{1g}$  mode from 404.2 cm<sup>-1</sup> for the monolayer MoS<sub>2</sub> to 406.7 cm<sup>-1</sup> for the MoS<sub>2</sub>-NSs, and this shift was attributed to the van der Waals interactions between neighboring NS layers $^{25,26}.$  $^{25,26}.$  $^{25,26}.$  Additionally, a red shift in the  $\mathrm{E_{\,2g}^1}$  mode was detected from 385.4 cm<sup>-1</sup> for the monolayer MoS<sub>2</sub> to 382.7 cm<sup>-1</sup> for the MoS2-NSs, resulting from stacking-induced structural changes, long-range Coulombic interlayer interactions<sup>25,27</sup>, or/and a slight tension owing to the bending deformation of the flake $28$ . The frequency shift behavior reflects structural changes after scrolling, and the changes were also observed for  $WS_{2}$ -,  $MoSe_{2}$ -, and  $WSe_{2}$ -NSs (Supplementary Fig. 12a–c).

Semiconducting TMD monolayers are direct-bandgap materials because of quantum confinement effects $29$ . The strong photoluminescence (PL) and a wide range of band gaps of the monolayer TMD semiconductors indicate their potential for use in 2D optoelectronics, such as optical energy conversion devices<sup>30</sup>or photodetectors $31$ . Despite having a stacking structure similar to that of bulk TMDs, these TMD-NSs were shown to have an obvious PL, in contrast with the bulk counterparts, in which PL is negligible<sup>29</sup>. Red luminescence from a typical  $MoS<sub>2</sub>-NS$  was clearly observed by fluorescence imaging, indicating that the PL was sufficiently strong to be imaged (Fig. 3b, c). In the PL spectra of Fig. 3d, A, B direct excitonic transitions<sup>[29](#page-6-0)</sup> from  $MoS_2$ -NSs appeared at around 1.80 eV (688 nm) and 1.95 eV (635 nm), matching their absorption resonances (Supplementary Fig. 13). Compared with monolayer  $MoS<sub>2</sub>$ , a red shift in the A exciton transition, approximately of

30–60 meV was found for  $MoS<sub>2</sub>-NSs$  and was accompanied by an increase in the full-width at half-maximum of 50–80 meV. The red shift of A peak, an indicator of the variation of the electronic structure, can be explained by the stacking and tension effects in the NSs<sup>29,[32,33](#page-6-0)</sup>, while PL peak broadening is presumably owing to the overlap of signals from several MoS<sub>2</sub> layers in a NS. Analogous PL behavior was also observed for the  $WS_2$ -, MoSe<sub>2</sub>-, and WSe<sub>2</sub>-NSs (Supplementary Fig. 12d–f).

It is well known that monolayer TMD atoms are fully exposed, which makes their optical and electrical properties sensitive to environmental factors. As shown in Fig. 3e, f, when exposed to  $NH<sub>3</sub>$ , the PL intensity of monolayer MoS<sub>2</sub> flakes was quenched rapidly and drastically; in contrast, the PL intensity of  $MoS<sub>2</sub>-NSs$ did not show a pronounced decrease. Apparently, the optical properties of  $MoS<sub>2</sub> - NSs$  were less dependent on the environment, which might be attributed to the self-encapsulated structure, where most of the  $MoS<sub>2</sub>$  flakes were rolled up inside and sealed by the outermost layers.

FET performance. Since TMD-NSs are semiconductors, and there are no previous reports on their FET properties, it is interesting to explore this issue here. Device characterization is easy for  $MoS<sub>2</sub>-NSs$  because the NSs have been produced on a  $SiO<sub>2</sub>/Si$  substrate. Defining source and drain electrodes were directly deposited on the top of  $MoS_2-NSs$  via a shadow mask. This process avoided organic impurity contamination that might occur during traditional electron-beam lithography. The FETs based on the  $MoS<sub>2</sub>-NSs$  exhibited good FET behavior, and the current/voltage characteristics conformed well to conventional transistor models in both saturated and linear regimes in  $N_2$ (Fig. [4](#page-4-0) and Supplementary Fig. 14). According to the equation below, we calculated the mobility in the saturation regime.

$$
I_{\rm D} = (W/2L)C_{\rm i}\mu(V_{\rm G} - V_{\rm T})^2,
$$

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Fig. 4 Electrical characteristics of MoS<sub>2</sub>-NSs (nanoscrolls). a Schematic representations of current conduction in transition metal dichalcogenide (TMD)nanoscrolls (NSs), TMDs monolayer, and TMDs multilayer under bias. The W represents the conduction channel width. The NSs made from the below TMDs monolayer have a much shorter W. For TMDs multilayer, the current only passes through the outermost shell layers. **b** SEM image of a typical MoS<sub>2</sub>-NS FET (scale bar, 10 μm). c Transfer characteristics of MoS<sub>2</sub>-NS tested in N<sub>2</sub> (red) and air (blue). d Output characteristics of MoS<sub>2</sub>-NS tested in N<sub>2</sub>. **e** Output characteristics of  $MoS<sub>2</sub>-NS$  tested in air

where L and W are the channel length and width, respectively,  $C_i$  is the insulator capacitance per unit area (11 nF cm<sup>-2</sup>),  $\mu$  is the FET mobility, and  $V_T$  is the extrapolated threshold voltage. More than 300 devices were tested. The electron mobility of NSs was measured to be mostly in the range 200–700 cm<sup>2</sup>  $\dot{V}^{-1} s^{-1}$ , with an on/off ratio of over  $1.0 \times 10^5$ . The mobility of MoS<sub>2</sub>-NSs, nearly 30-fold greater than that of monolayer  $MoS<sub>2</sub>$  flakes (10–20 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>)<sup>[34](#page-6-0)</sup>, should be ascribed to the high current density for compact scroll topology. Invoking an Archimedean spiral model, we quantified the change in the conduction channel widths upon scroll (the detailed analysis is provided in Supplementary Note [3](#page-3-0)). A considerable reduction in conduction channel widths was achieved after scrolling, while there is no significant decline in the current for the flakes upon scrolling process (Supplementary Fig. 15, 16). This led to a higher current density and mobility. In addition, benefiting from the compact scrolling, conductive layers in a NS compactly stacked layer by layer. The adjacent layers acted as ideal substrates for each other with atomically smooth surface, relatively no dangling bonds and charge traps, which also enhanced electron transporting efficiency<sup>[35](#page-6-0)</sup>. With the similar stacked layer structures, however, the  $MOS_2$ -NSs also exhibited much greater electron mobility than did the multilayer  $MoS<sub>2</sub>$ flakes<sup>[36](#page-6-0)</sup>. It can be explained that in a  $MoS<sub>2</sub>-NS$ , electrons are transported through the whole scrolled  $MoS<sub>2</sub>$  flake, while in multilayer  $MoS<sub>2</sub>$  flakes, electrons are blocked by van der Waals gaps while migrating to inner layers; thus, only a few shell layers contribute to the carriers transmission (Fig. 4a). The monolayer  $MoS<sub>2</sub>$ -based FETs were observed to be very sensitive to atmosphere. However, there was no obvious decrease in FET mobility for the  $MoS_2-NSs$  when they were tested in air (Fig. 4c-e)

because most of each electrically active  $MoS<sub>2</sub>$  layers is encapsulated within each NS and the compact stacked (with an interlayer distance of 0.62 nm) morphology can efficiently prevent oxygen from intercalating into the MoS<sub>2</sub>-NSs to deteriorate transistors performance. The high mobility and unique self-encapsulated morphology make  $MoS<sub>2</sub>$ -NSs promising in FETs.

Hybrid TMD-NSs. As mentioned above, TMD-NSs are inert to the outsides owing to their self-encapsulated structures; however, they are inclined to be doped and form composites because of their internal open topological structures. With a pretreatment, TMD flakes can be hybridized with diverse loads, which become sandwiched into the van der Waals gaps upon flakes scrolling (Fig. [5](#page-5-0)a). In this paper, we produced MoS<sub>2</sub>-NSs hybridized with gold nanoparticles (AuNPs) (Fig. [5b](#page-5-0)), graphene oxide (GO) (Fig. [5](#page-5-0)c), pentacene (Fig. [5](#page-5-0)d), copper (II) phthalocyanine (CuPc) (Fig. [5e](#page-5-0)), poly{2,2ʹ-[(2,5-bis (2-hexyldecyl) −3,6-dioxo- −2,3,5,6-tetrahydropyrrolo[3,4-c] pyrrole-1,4-diyl)dithiophene] −5,5ʹ-diyl-alt-thiophen- −2,5-diyl} (PDPP3T) (Fig. [5](#page-5-0)f), DNA (Fig. [5](#page-5-0)g) and polypeptide (Fig. [5](#page-5-0)h). In contrast with nanotubes, the lattice spacing of NSs can adjust automatically upon hybridization with foreign substances of different sizes, and this adjustment can be evaluated by HR-TEM. Figure [5b](#page-5-0)–h displayed the lattice spacing of hybrid NSs, where lattice expansion was observed for all composite NSs, and some composite components, such as AuNPs, were large enough to be directly visualized (Fig. [5b](#page-5-0)). Moreover, Raman spectroscopy was utilized to verify the presence of foreign substances. As a typical example, the intercalation of GO into NSs was verified by the emergence of the D, G, and 2D peaks of GO (around 1350, 1600, 2671 cm<sup>-1</sup>)<sup>37</sup>, as well as vibrational modes (around 383, 406 cm<sup>-1</sup>) for  $MoS<sub>2</sub>$  (Supplementary Fig. 17a). In addition, TMD-NSs can be conferred with additional properties and functions upon hybridization with other functional materials, such as organic small molecules, polymers, nanoparticles, and 2D materials, as well as biological substances. For example, after hybridization with PDPP3T, a new fluorescence peak appeared at about 1.42 eV (873 nm) for  $MoS<sub>2</sub>-NSs$  (Supplementary Fig. 17b), indicating that the hybridization broadened  $MoS<sub>2</sub>-NS$  emission spectrum and could eventually extend their optical applications. Furthermore, the intercalated NS by PDPP3T exhibited a good stability in air, due to the self-encapsulated structure (Supplementary Fig. 18).

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Fig. 5 Hybrid transition metal dichalcogenide (TMD)-nanoscrolls (NSs). a Schematic showing the cross section and sidewalls of hybrid TMD-NSs. **b-h** HR-TEM images of sidewalls of MoS<sub>2</sub>-NSs hybridized with AuNPs (b), GO (c), pentacene (d),  $CuPc$  (e), PDPP3T (f), DNA (g) and polypeptide (h). AuNPs approximately 5 nm in diameter were clearly observed between the walls; the GO flakes were also distinguishable, as shown in c. Lattice expansion was displayed in all hybrid NSs above. The bright white walls represent the MoS<sub>2</sub> layers in  $(c-h)$  (scale bars, 10 nm in **b** and  $2 \text{ nm}$  in  $c-h$ )

#### **Discussion**

In summary, we have explored a simple method for producing high-quality, tightly scrolled TMD-NSs from CVD-grown monolayer TMD flakes with the assistance of one drop of an ethanol solution. An obvious PL was shown in TMD-NSs, with a red shift in the direct excitonic transitions. The electron mobility of the  $MoS<sub>2</sub>-NSs$  was approximately 30-fold greater than that of the monolayer  $MoS<sub>2</sub>$  flakes. This scrolling strategy will be a candidate for increasing mobility of other 2D materials. Owing to their unique self-encapsulated structures, these TMD-NSs show stable optical and electronic properties in different atmospheres. The HR-TEM results show the internal open topological structure of the TMD-NSs, which serve as effective carriers for accommodating external substance of different sizes in their tunable van der Waals gaps. The hybridization with other functional materials will confer TMD-NSs with additional properties and functions attractive for potential application.

#### Methods

For MoS<sub>2</sub> monolayers, the growth substrate was placed face-down above a ceramic boat containing MoO<sub>3</sub> powder (Sigma Aldrich, 99.998%, 20 mg). The ceramic boat was then loaded in the heating zone center of the furnace tube, where another boat containing 80 mg of sulfur (Alfa Aesar, 99.999%) was located upstream. The tube was first purged with ultrahigh-purity argon for 10 min at a flow rate of 200 sccm. Then, the furnace was heated from room temperature to 500 °C in 12 min, 500–720 °C in another 20 min, and then stayed at 720 °C for 5 min. The sulfur was heated to 130 °C with a separate heat setup as the furnace reached 720 °C. Finally, the temperature was cooled from 720 to 570 °C in 20 min before opening the furnace for rapid cooling. Argon as carrier gas was maintained at a flow rate of 10 sccm throughout the growth process.

The growth of  $WS_2$  monolayers was similar to that of  $MoS_2$ , except that  $WO_3$ (Alfa Aesar, 99.998%, 10 mg) was used as one of the precursors and the highest growth temperature was set at 825 °C, with sulfur at 150 °C.

 $WSe<sub>2</sub>$  monolayers were grown on substrates with  $WO<sub>3</sub>$  (10 mg) and selenium powders (Alfa Aesar, 99.999 + %, 1–5 mm) as precursors. The growth recipe was as follows: ramp the furnace from room temperature to 825 °C in 20 min, sit 20 min at 825 °C, cool to 600 °C in 20 min, and open the furnace for rapid cooling. The selenium powders were heated to 300 °C when the furnace temperature rose to 825 °C and were maintained at this temperature for 40 min. 1.5 sccm hydrogen and 23.5 sccm argon were used during the growth process.

 $MoSe<sub>2</sub> monolayers were grown via a procedure similar to that used for  $WSe<sub>2</sub>$$ monolayers, except that the source precursors were  $MoO<sub>3</sub>$  and selenium, and the selenium was kept at 400 °C during the growth process.

Fabrication of TMD-NSs. TMD-NSs were fabricated by dropping an ethanol aqueous solution onto the CVD-grown TMD monolayers and allowing it to dry naturally. Rod-like NSs were then obtained on the substrate. The concentration of the ethanol aqueous solution ranged from 5:1 to 1:1 (volume ratio of ethanol: water), with an optimum ratio of 2:1.

 $MoS<sub>2</sub>-NSs$  arrays. A  $MoS<sub>2</sub>$  polycrystalline film was first patterned into parallel ribbons with defined width and direction using a FIB. A drop of pure ethanol was dropped on the patterned film, and then a coverslip was put on top of it. A thin liquid layer of pure ethanol can be formed between the film and the coverslip. An ethanol solution (1:1) was then dropped on the edge of the coverslip. Drove by the concentration gradient, the water in the solution will slowly infiltrate into the pure ethanol under the coverslip in the gradient direction (as shown in Supplementary Fig. 8). The diffusion rate can be adjusted by the concentration of the ethanol solution. In this way one side of a ribbon was first controlled to scroll to produce a long NS. The long NS array was then fabricated and further arranged into periodic arrays with a controlled length by FIB etching a second time.

Hybrid MoS<sub>2</sub>-NSs fabrication. Pentacene/MoS<sub>2</sub>-NSs: a 1-nm-thick layer of pentacene was deposited on the MoS<sub>2</sub> monolayers on a  $SiO<sub>2</sub>/Si$  substrate by vacuum thermal evaporation, followed by immersion into a modulated ethanol solution (5:1). After 10 min, the substrate was picked up and allowed to dry naturally. The pentacene/MoS<sub>2</sub>-NSs were obtained. The CuPc/MoS<sub>2</sub>-NSs were prepared in the same way.

PDPP3T/MoS2-NSs: A PDPP3T solution was prepared by dissolving 2 mg of PDPP3T powder in 1 ml of toluene. After stirring overnight, the PDPP3T solution was spin-coated onto the  $MoS<sub>2</sub>$  monolayers at a speed of 5000 rpm. After annealing at 70 $^{\circ}$ C for 20 min, the PDPP3T-coated MoS<sub>2</sub> monolayers were then dipped into an ethanol solution (5:1). After 10 min, the  $MoS<sub>2</sub>$  monolayers were picked up and dried naturally.

 $GO/MoS<sub>2</sub>$ -NSs: This procedure is similar to that used for PDPP3T/MoS<sub>2</sub>-NSs. Freshly exfoliated GO sheets were dispersed in n-hexane and then spin-coated onto MoS2 monolayers. After dropping an ethanol solution (2:1) onto the samples, the GO/MoS<sub>2</sub>-NSs were obtained.

Au/MoS2-NSs: MoS2 monolayers were dipped into an ethanol solution (2:1) containing AuNPs. Ten minutes later, the  $MoS<sub>2</sub>$  monolayers were picked up and dried naturally. Similar procedures were used for producing both  $\overline{DNA}/\overline{MoS}_2$ - and polypetide/MoS<sub>2</sub>-NSs.

 $MoS<sub>2</sub>-NS-based FET$ . Freshly made monolayer  $MoS<sub>2</sub>$  flakes on a SiO<sub>2</sub>/Si substrate were kept in nitrogen for 48 h. Then the  $MoS<sub>2</sub>$  flakes were scrolled into NSs with one drop of ethanol solution (2:1) in nitrogen atmosphere. Bottom-gated transistors based on these  $MoS_2-NSs$  were fabricated on the  $SiO_2/Si$  substrate without additional transfer. 100 nm of Au was deposited on the top of  $MoS<sub>2</sub>$ -NSs as source and drain electrodes with a shadow mask by vacuum thermal evaporation at a rate of about 0.3 Å s<sup>−</sup>1. The electrical measurements were conducted under a nitrogen atmosphere and ambient condition, respectively.

Characterization. Raman and PL measurements were performed with a 532 nm laser under ambient conditions (inVia-Reflex). All optical images and the video were captured with a Nikon Eclipse LV100D. The fluorescence microscope characterization was performed using a fluorescence microscope equipped with a mercury lamp as the excitation light source. Large-area  $\mathrm{MoS}_{2}$  films and long MoS2-NSs were etched via FIB-SEM (FEI, Helios Nanolab G3 CX). SEM

CVD growth of TMD monolayers. TMD monolayers were grown by an atmospheric pressure CVD. The growth substrates, including  $SiO<sub>2</sub>/Si$ ,  $Si<sub>3</sub>N<sub>4</sub>$  and sapphire, were cleaned in acetone, isopropanol, and deionized water and then dried under high-purity nitrogen flow.

<span id="page-6-0"></span>(Hitachi S-4800) and HR-TEM (JEOL JEM-2100F) were employed to image the samples. Current–voltage curves were obtained by a three-probe station under a nitrogen atmosphere and ambient conditions (Keithley 4200 SCS).

Data availability. The data that support the findings of this study are available from the corresponding author on reasonable request.

Received: 24 August 2017 Accepted: 8 March 2018 Published online: 03 April 2018

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### Acknowledgements

We thank Bo Guan and Xiang Li for the assistance with FIB-SEM. We thank Bo Guan for the assistance in SAED analysis. This work was supported by the Strategic Priority Research Programme of the National Natural Science Foundation (21573253), and the Chinese Academy of Sciences (XDB12010000).

### Author contributions

X.C. performed the experiments, analyzed the data, and wrote the paper. J.Z. designed the work, performed the experiments, conceptualized the work, and analyzed the data. Z.K., D.H., Y.H., H.S., C.D., and D.Z. performed some characterization experiments and analyzed data. E.G. and Z.X. performed the theoretical analysis.

### Additional information

Supplementary Information accompanies this paper at [https://doi.org/10.1038/s41467-](https://doi.org/10.1038/s41467-018-03752-5) [018-03752-5.](https://doi.org/10.1038/s41467-018-03752-5)

Competing interests: The authors declare no competing interests.

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## Supplementary Information

# Rolling up Transition Metal Dichalcogenide Nanoscrolls via One Drop of Ethanol

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## Supplementary Figures



Supplementary Figure 1. a, Optical image of built-in tension released monolayer  $MoS<sub>2</sub>$  after transferred from growth substrate to another bare substrate. **b**, Optical image of the transferred sample after immersing into the ethanol solution. Almost no scroll was observed in the built-in tension released sample. c, Optical image of the MoS<sub>2</sub> sample after sharp intercalation of pure water. Loosely curled and disordered nanoscrolls with large unscrolled film were observed. (Scale bars, 200 μm in a, b and 20 μm in c).



Supplementary Figure 2.  $MoS<sub>2</sub>-NSs$  made in the aqueous solution of methanol (a), tetrahydrofuran (b), dimethylformamide (c), and N-methyl-2-pyrrolidone (d). (Scale bars, 100 μm in a and 50 μm in b-d). Tight, straight scrolling with high yield can also be obtained with different solutions at the optimized ratio.



Supplementary Figure 3.  $MoS<sub>2</sub>-NSs$  made from monolayer  $MoS<sub>2</sub>$  flakes grown on non-silicon substrates. a, Optical image of  $MoS<sub>2</sub>$  monolayers grown on a  $Si<sub>3</sub>N<sub>4</sub>$ substrate. The dark yellow represents  $MoS<sub>2</sub>$  monolayers, and the light yellow is the  $Si<sub>3</sub>N<sub>4</sub>$  substrate. **b**, Optical image of MoS<sub>2</sub>-NSs made from MoS<sub>2</sub> monolayers grown on the  $Si<sub>3</sub>N<sub>4</sub>$  substrate. c, Optical image of  $MoS<sub>2</sub>$  monolayers grown on a sapphire substrate. The white area represents  $MoS<sub>2</sub>$  monolayers, and the brown area represents the sapphire substrate. d, Optical image of  $MoS<sub>2</sub>-NSs$  rolled up from  $MoS<sub>2</sub>$ monolayers grown on the sapphire substrate.



Supplementary Figure 4. The evolution in contrast color of  $MoS<sub>2</sub>$  flakes when treated by ethanol solution. a, Optical image of  $MoS<sub>2</sub>$  flakes on a  $SiO<sub>2</sub>/Si$  substrate before dropping an ethanol solution.  $\mathbf{b}$ , Optical image of MoS<sub>2</sub> flakes just covered with the ethanol solution. c, Optical image of  $MoS<sub>2</sub>$  flakes when a small amount of ethanol solution intercalating into  $MoS<sub>2</sub>$  flakes and the substrate  $d$ , Optical image of  $MoS<sub>2</sub>$  flakes when a large amount of ethanol solution intercalating into  $MoS<sub>2</sub>$  flakes and the substrate. (Scale bars, 100 μm).



Supplementary Figure 5. a, The TEM image of a typical  $MoS<sub>2</sub>-NS$ . b, a schematic illustration of the  $MoS_2-NS$  in (a). c, Potential energy analysis of a flat  $MoS_2$  sheet (blue line), a flat  $MoS<sub>2</sub>$  sheet with solvent intercalation at its interface with the substrate (red dash line), and a scrolled  $MoS<sub>2</sub>$  sheet (red line), plotted as the function of length l.



Supplementary Figure 6. SEM image of an aged  $MoS<sub>2</sub>$  sample on a  $SiO<sub>2</sub>/Si$ substrate.



Supplementary Figure 7. Chirality of the monolayer MoS<sub>2</sub> single crystal scrolling. **a**, SEM image of a MoS<sub>2</sub>-NS rolled up from a single-crystal MoS<sub>2</sub> monolayer. **b**, High-magnification SEM image of the rectangular area in (a). c-e, Schematic of scrolling chirality of TMD-NSs made from single-crystal flakes. The single-crystal TMDs flakes curled along the zigzag edge. (Scale bars, 10 μm in a and 500 nm in b)



Supplementary Figure 8. a, Schematic of the preparation of a  $MoS<sub>2</sub>-NS$  array. b, A long  $MoS_2-NS$  array with a maximum area of about 3 mm×3 mm on a  $SiO_2/Si$ substrate.



Supplementary Figure 9. Infrared (IR) and Raman analysis of  $MoS<sub>2</sub>-NSs.$  a, IR spectrum of ethanol (red) and  $MoS<sub>2</sub>-NSs$  (blue). **b**, Raman spectra of ethanol (red) and  $MoS<sub>2</sub>-NSs$  (blue). A weak IR band at 481 cm<sup>-1</sup> was assigned to Mo-S vibration. In the Raman spectra, vibration modes (at  $382.9 \text{ cm}^{-1}$ ,  $407.3 \text{ cm}^{-1}$ ) were clearly observed for  $MoS_2-NSS$ , without modes for ethanol in the NS. IR and Raman spectra both showed there was no residual ethanol in the nanoscrolls. The IR analysis was carried out under ambient conditions on a microscope setup (Nicolet iN10-iZ10). Raman characterization was performed with a 785 nm laser (inVia-Reflex).



Supplementary Figure 10. a, SEM image of a  $MoS_2-NS$ . b, An enlarged image of the marked area in (a).



Supplementary Figure 11. TEM image and the selected-area diffraction (SAED) pattern (inset) of an individual  $MoS_2-NS$ . Two sharp diffraction spots, indexed to (002) plane, indicated the uniformity of the interlayer spacing in the NS. According to Bragg's diffraction equation, the interlayer spacing of the NS was about 6.18 Å, in good agreement with the HR-TEM observation.



Supplementary Figure 12. Optical characterization of  $WS_2$ -,  $MoSe_2$ -, and WSe<sub>2</sub>-NSs. a-c, Raman spectra of CVD-grown  $WS_2$ , MoSe<sub>2</sub>, and WSe<sub>2</sub> monolayers (ML, black) and NSs (red), respectively.  $d$ -f, PL spectra of CVD-grown WS<sub>2</sub>, MoSe<sub>2</sub>, and WSe<sub>2</sub> monolayers (black) and NSs (red) (532 nm excitation).



Supplementary Figure 13. Optical absorption characterization of  $MoS<sub>2</sub>$  monolayers ( $MoS<sub>2</sub>-ML$ ) and  $MoS<sub>2</sub>-NSs$ . Monolayer  $MoS<sub>2</sub>$  and the thin film of NSs were both measured on a fused quartz substrate in an integrating sphere set-up (UV-2600). The NS film was prepared by the accumulation of the multiple transferred nanoscrolls from  $SiO<sub>2</sub>/Si$  substrate and the spectra were normalized for better comparison. Obvious red shift in absorption resonances was observed after scrolling, corresponding to the red shift of exciton transitions in the PL spectra.



Supplementary Figure 14. High-resolution scanning electron microscope (HR-SEM) and atomic force microscope (AFM) characterization of  $MoS<sub>2</sub>-NS$ FET in Figure 4b. a, HR-SEM image of the  $MoS<sub>2</sub>-NS$  FET. b, AFM image of the  $MoS<sub>2</sub>-NS FET showed in (a). c, HR-SEM image of part of the MoS<sub>2</sub>-NS FET in (a). d,$ AFM image of the  $MoS_2-NS FET$  showed in (c). (Scale bars, 5  $\mu$ m in a,b and 2 $\mu$ m in c,d)



Supplementary Figure 15. The schematic of a monolayer TMD flake rolled into a NS along the conduction channel width direction. W1, W2 represent the conduction channel widths of TMD flake and the NS respectively.  $\theta$  is the torsion angle expressed in polar coordinates.



Supplementary Figure 16. Electrical characterizations of monolayer  $MoS<sub>2</sub> FETs.$ a, Transfer curve of a back-gated monolayer  $MoS<sub>2</sub> FET$  measured in  $N<sub>2</sub>$  at room temperature. The source-drain bias was fixed at 60 V. Inset: Photograph of the FET device made from a monolayer  $MoS_2$  flake grown on a  $SiO_2/Si$  substrate by CVD. The monolayer  $MoS<sub>2</sub>$  flake has similar dimensions as the flake that rolled into the NS in Figure 4b. b, Output curve of the monolayer  $MoS<sub>2</sub> FET. Experimented results show$ that the flake FET has a similar saturated current to that of the NS FET under the same bias. There is no significant decline in saturated current for the flakes upon scrolling process.



Supplementary Figure 17. Optical characterization of hybrid  $MoS<sub>2</sub>-NSs$ . a, Raman spectrum of  $GO/MoS_2$ -NSs. The inset is a magnified view of the range of 370-420 cm<sup>-1</sup>. **b**, PL spectra of MoS<sub>2</sub>-NSs (black) and PDPP3T/MoS<sub>2</sub>-NSs (red) (532 nm excitation).



Supplementary Figure 18. a, HR-TEM images of PDPP3T/ $MoS_2$ -NSs one month after the first intercalation. b, Raman spectra of PDPP3T/MoS<sub>2</sub>-NSs placed in ambient condition for 0 day (black line) and one month (red line). c, PL spectra of PDPP3T/MoS<sub>2</sub>-NSs placed in ambient condition for 0 day (black line) and one month (red line). (532 nm excitation)

### Supplementary Notes

# Supplementary Note 1. The role of immersing solution in CVD-based TMD flake scrolling

To clarify the role of the immersing solution in TMD flake scrolling, an analog experiment was carried out. CVD-grown  $MoS<sub>2</sub>$  flakes were transferred from growth substrate to another bare  $SiO<sub>2</sub>/Si$  substrate with a polymer-based method. In this process, the built-in tension in the flakes can be released because of the elastic polymer support [1]. Then the transferred sample was immersed into the ethanol solution for 30 minutes. Almost no scrolls were observed in the transferred sample after the procedures above, as shown in Supplementary Fig. 1a and b. The result indicated that the tension to scroll the flake was not induced by the immersing solution. It was observed that the immersing solution here was to intercalate into the flake and the substrate, thereby weakening the adhesion between them. This was crucial for built-in tension to initiate the flake scrolling, as analyzed in theory section (Supplementary Note 2). The existence of water was found to be indispensable to the intercalation. A strong interaction of hydrogen bond can be formed between water and the growth substrate  $(SiO_2/Si, Si_3N_4$  or sapphire), which is stronger than the Van der Waals' force between TMD flakes (hydrophobic) and the substrate. Thus after being dropped on the sample, the water tends to substitute the flakes to contact with the substrate, leading to the intercalation between flakes and the substrate. Instantaneous intercalation happened to the samples when only water used. But the sharp intercalation often caused large pieces of TMD film peeled off together from the substrate and floated in the water, giving rise to loosely curled, disordered nanoscrolls with large film unscrolled (Supplementary Fig. 1c). A moderate intercalation speed is required for the formation of high-quality nanoscrolls. Under a low intercalation speed, the substrate can act as a template to enable the flake to scroll straightly and tightly. Mixing water with water-soluble organic solvents, such as ethanol, methanol and dimethylformamide, was observed to slow down the intercalation. When only pure organic solvent was used, the intercalation proceeded quite slowly, and even cannot be observed. Optimizing the ratio of water and organic solvents in solution, high-quality scrolling of TMD flakes can be achieved. With the same concentration of aqueous solution, the scrolling speed slightly increased with the polarity of the organic solvent. High-quality scrolling can also be obtained with different solutions at the optimized ratio, as shown in Supplementary Fig. 2.

### Supplementary Note 2. Energy analysis of the formation process of a  $MoS<sub>2</sub>-NS$

To understand the formation mechanism of a  $MoS<sub>2</sub>-NS$  as a result of the release of elastic energy upon solvent intercalation, we present a comparative energy analysis of flat and scrolled  $MoS_2$  sheets on a  $SiO_2/Si$  substrate.

For a stress-free  $MoS_2$  sheet with length of l and width w adhered on a  $SiO_2/Si$ substrate, the potential energy composes of the interfacial adhesion energy  $U_A = -\gamma_1 A$ , where  $\gamma_1$  is the areal adhesion energy density and  $A = w l$  is the contact area between  $MoS<sub>2</sub>$  sheet and substrate, as well as the in-plane tensile energy of the  $MoS<sub>2</sub>$  sheet resulted from the mismatch in thermal strain as the  $MoS<sub>2</sub>/substrate$  system is cooled from 720-825 °C to the room temperature, which is  $U_I = E_{2D}[(\alpha_1 - \alpha_2)\Delta T]^2 w l/2$ . Here  $E_{2D}$  is the in-plane tensile stiffness of monolayer MoS<sub>2</sub> sheet,  $\alpha_1$  and  $\alpha_2$  are the thermal expansion coefficients of MoS<sub>2</sub> and substrate, respectively, and  $\Delta T$  is the change of temperature. The total potential energy  $(U_f)$  of the hybrid with a flat MoS<sub>2</sub> sheet on  $SiO<sub>2</sub>/Si$  substrate is then

$$
U_{\rm f} = E_{2D}[(\alpha_1 - \alpha_2)\Delta T]^2 w l/2 - \gamma_1 w l \tag{1}
$$

For a scrolled  $MoS_2$  sheet on the  $SiO_2/Si$  substrate, the total energy (U) includes the self-adhesion energy of the MoS<sub>2</sub> sheet  $U_s = -\gamma_2(l - l_s)$ , and the out-of-plane bending energy  $(U_B)$  of MoS<sub>2</sub>,

$$
U_{\rm B} = \frac{\pi D}{h} \ln \left( \frac{\sqrt{lh/\pi + r_0^2}}{r_0} \right),
$$

where  $\gamma_2$  is interlayer adhesion energy density between MoS<sub>2</sub> sheet itself,  $l_s = 2\pi r_0$  is the critical length of cylindrical nucleus of the scroll,  $h$  is the interlayer spacing of  $MoS<sub>2</sub>-NSs$ , and  $r<sub>0</sub>$  is the radius of the cylindrical nucleus, as shown in Supplementary

Fig. 5a, b. The total potential energy of the hybrid with  $MoS<sub>2</sub>-NS$  on  $SiO<sub>2</sub>/Si$  substrate is then

$$
U = \begin{cases} \frac{\pi D w}{h} \ln \left( \frac{\sqrt{l h / \pi + r_0^2}}{r_0} \right) & (l < 2\pi r_0) \\ \frac{\pi D w}{h} \ln \left( \frac{\sqrt{l h / \pi + r_0^2}}{r_0} \right) - \gamma_2 w (l - 2\pi r_0) & (l \ge 2\pi r_0) \end{cases}
$$
(2)

Using the typical measured values of parameters in our previous analysis, that are  $\Delta T$ about 750 K,  $r_0$  about 10 nm (from our experimental setup),  $h = 0.621$  nm,  $y_1 = 0.170$ J m<sup>-2</sup> [2],  $\gamma_2$  = 0.140 J m<sup>-2</sup> [3],  $E_{2D}$  = 122.30 N m<sup>-1</sup> [4], D = 6.29 eV [4],  $\alpha_1 = 7 \times 10^{-6} \text{ K}^{-1}$ [5] and  $\alpha_2 = 2.6 \times 10^{-6}$  K<sup>-1</sup> [6]. For a flat sheet to bend into a scroll, a cylindrical nucleus has to be formed with a critical length of nucleation  $l_s = 2\pi r_0$  (around 62.8 nm) as observed in our experiments, and an energy barrier  $E_b = 1.9 \text{ eV nm}^{-1}$ , as calculated from supplementary equation 2.

As the solvent intercalation proceeds, the interlayer spacing between  $MoS<sub>2</sub>$  and substrate is increases and the adhesion energy  $\gamma_1$  will decrease abruptly and we assume that  $\gamma_1$  become zero (about 0 J m<sup>-2</sup>) after the delamination.

We plot the total potential energy per width for a flat  $MoS<sub>2</sub>$  flake (blue line), a flat  $MoS<sub>2</sub>$  flake with intercalated solvent (red dash line), and a scrolled  $MoS<sub>2</sub>$  (red line), as a function of the contour length of sheets  $l$  (Supplementary Fig. 5c). Our energy analysis shows that there is a relatively low energy barrier of 0.3 nJ  $m^{-1}$  or 1.9 eV nm<sup>-1</sup> for the scroll to nucleate below the critical length for scroll nucleation  $l_s$ , compared with the flat  $MoS<sub>2</sub>$  with intercalated solvent. As shown in the inset of Supplementary Fig. 5c, as the contour length of  $MoS<sub>2</sub>$  sheets exceeds 0.45 µm, the potential energy of a flat  $MoS<sub>2</sub>$  with solvent intercalation (red dash line) is higher than the energy barrier to form scrolled MoS<sub>2</sub>. This result indicates that the strain energy induced from the thermal mismatch of a flat  $MoS<sub>2</sub>$  is high enough to activate the formation of scrolls after solvent intercalation and decreasing the adhesion energy with substrate (Supplementary Fig. 5c). It should be remarked here that the conclusion holds on the condition that  $\gamma_1$  becomes lower than  $\gamma_2$  after solvent intercalation.

### Supplementary Note 3. Analysis of the effect of compact scroll on FET mobility

We model the cross-section of a TMD-NS as an Archimedean spiral [7], as shown in Supplementary Fig. 15. The change of the scroll radius  $r$  along the length depends on the torsion angle  $\theta$  expressed in polar coordinates:

$$
r = a\theta + b \tag{1}
$$

where  $a=t/2\pi$  is the spiral constant; t is the radius change per one full turn, approximately equal to the interlayer distance in the NS; b is the scroll radius at  $\theta=0$ , i.e., the inner radius of the NS. As determined from TEM images, the  $t$  is approximately 0.6 nm and b is in the range of 5-20 nm for TMDs-NSs.

The arc length s of the Archimedean spiral is an integral by  $\theta$ .

$$
s = \int_0^s ds = \int_0^{\theta} \sqrt{(dr)^2 + (r d\theta)^2}
$$
 (2)

Substitution of supplementary equation (1) into supplementary equation (2) yields

$$
s = \int_0^{\theta} \sqrt{a^2 + (a\theta + b)^2} \, d\theta \tag{3}
$$

Based on the integral formula, the supplementary equation (3) can be computed as

$$
s = \frac{a\theta + b}{2a} \sqrt{a^2 + (a\theta + b)^2} + \frac{a}{2} \ln \left[ b + a\theta + \sqrt{a^2 + (a\theta + b)^2} \right] - \frac{b}{2a} \sqrt{a^2 + b^2}
$$

$$
- \frac{a}{2} \ln \left( b + \sqrt{a^2 + b^2} \right) \tag{4}
$$

As shown in Supplementary Fig. 15, a TMD flake scrolled into a NS along one edge. The conduction channel width  $W1$  for the TMD flake is just the arc length  $s$  of the NS spiral, while the conduction channel width  $W2$  for the NS is equal to its diameter  $2r$ . Thus, the ratio of the conduction channel widths before and after scrolling can be described as

$$
\frac{W1}{W2} = \frac{s}{2r} \tag{5}
$$

Substitution of supplementary equations (1) and (4) into supplementary equation (2)

yields

$$
\frac{W1}{W2} = \frac{1}{4a} \sqrt{a^2 + (a\theta + b)^2} + \frac{a}{4(a\theta + b)} \ln \left[ b + a\theta + \sqrt{a^2 + (a\theta + b)^2} \right]
$$

$$
-\frac{b}{4a(a\theta + b)} \sqrt{a^2 + b^2} - \frac{a}{4(a\theta + b)} \ln \left( b + \sqrt{a^2 + b^2} \right) \tag{6}
$$

Given that *a* is relatively small value of  $\sim 0.095$ , the supplementary equation (6) is simplified to

$$
\frac{W1}{W2} = \frac{\theta}{4} - \frac{b^2}{4a(a\theta + b)} + \frac{b}{4a}
$$
 (7)

Assuming that the flake rolled by full turns with  $\theta = 2n\pi$  (*n*, the scroll turns, i.e., layer number in the NS), the dependence of the ratio of the conduction channel widths on the scroll turns can be expressed as follows:

$$
\frac{W1}{W2} = \frac{n\pi}{2} - \frac{b^2}{4a(2an\pi + b)} + \frac{b}{4a}
$$
 (8)

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